Rotational KE: is the sum of translational KE of all individual particles
 Let the object consists of N particles of mass $\mathrm{m}_{1}, \mathrm{~m}_{2}, \ldots \mathrm{~m}_{\mathrm{N}}$ at respective perpendicular distance of $r_{1}, r_{2}, \ldots, r_{N}$ from the axis of rotation. As the rigid body is rotating, all these particles are performing UCM with constant angular speed $\omega$ about an axis perpendicular to the plane of paper. But they all have different linear speeds $v_{1}=r_{1} \omega, v_{2}=r_{2} \omega \ldots . ., v_{N}=r_{N} \omega$
Translational $K E$ is $K E_{1}=1 / 2 m_{1} v_{1}{ }^{2}=1 / 2 m_{1} r_{1}{ }^{2} \omega^{2}, K E_{2}=1 / 2 m_{2} v_{2}{ }^{2}=1 / 2 m_{2} r_{2}{ }^{2} \omega^{2}, \ldots$. Rotational $K E=1 / 2 m_{1} r_{1}{ }^{2} \omega^{2}+1 / 2 m_{2} r_{2}{ }^{2} \omega^{2}+\ldots \ldots . .+1 / 2 m_{N} r_{N}{ }^{2} \omega^{2}$
$=1 / 2\left(m_{1} r_{1}{ }^{2}+m_{2} r_{2}{ }^{2}+\ldots .+m_{N} r_{N}{ }^{2}\right) \omega=1 / 2 I \omega^{2}$
where $\mathrm{I}=\mathrm{m}_{1} \mathrm{r}_{1}{ }^{2}+\mathrm{m}_{2} \mathrm{r}_{2}{ }^{2}+\ldots .+\mathrm{m}_{N} \mathrm{rr}^{2}=\sum_{i=1}^{N} m_{i} r_{i}{ }^{2}$ is the rotational inertia or moment of Inertia (MI) of the object about the given axis of rotation. NOTE: MI depends on individual masses and the distribution of these masses about the axis of rotation.

## Angular Momentum in terms of M1:

Let the object consists of $N$ particles of mass $m_{1}, m_{2}, \ldots m_{N}$ at respective perpendicular distance of $r_{1}, r_{2}, \ldots, r_{N}$ from the axis of rotation. As the rigid body is rotating, all these particles are performing UCM with constant angular speed $\omega$ about an axis perpendicular to the plane of paper. But they all have different linear speeds $v_{1}=r_{1} \omega, v_{2}=r_{2} \omega \ldots \ldots, v_{N}=r_{N} \omega$. These velocities are along the tangent.
Linear momentum is $p_{1}=m_{1} v_{1}=m_{1} r_{1} \omega, p_{2}=m_{2} v_{2}=m_{2} r_{2} \omega, \ldots \ldots .$.
Angular momentum is $L_{1}=m_{1} r_{1}{ }^{2} \omega, L_{2}=m_{2} r_{2}{ }^{2} \omega, \ldots . ., L_{N}=m_{N} r^{2} \omega$ Since all have same direction, the magnitude of angular momentum is $\mathrm{L}=\mathrm{m}_{1} \mathrm{r}_{1}{ }^{2} \omega+\mathrm{m}_{2} r_{2}{ }^{2} \omega+\ldots . .+\mathrm{m}_{N} r_{N}{ }^{2} \omega=\left(\mathrm{m}_{1} r_{1}{ }^{2}+\mathrm{m}_{2} r_{2}{ }^{2}+\ldots . .+\mathrm{m}_{N} r_{N}{ }^{2}\right) \omega=\mathrm{I} \omega$ where $I=m_{1} r_{1}^{2}+\mathrm{m}_{2} r_{2}{ }^{2}+\ldots .+\mathrm{mNr}^{2}=\sum_{i=1}^{N} m_{i} r_{i}^{2}$ is the moment of Inertia (MI) of the object about the given axis of rotation.

NOTE: $\mathrm{L}=\mathrm{I} \omega$ is analogous to linear momentum $\mathrm{p}=\mathrm{mv}$, if moment of Inertia (I) replaces mass, which is its physical significance.

## Torque in terms of Ml:



Let the object consists of $N$ particles of mass $\mathrm{m}_{1}, \mathrm{~m}_{2}, \ldots \mathrm{~m}_{\mathrm{N}}$ at respective perpendicular distance of $r_{1}, r_{2}, \ldots, r_{N}$ from the axis of rotation. As the rigid body is rotating, all these particles are performing CM with constant angular acceleration $\alpha$ about an axis perpendicular to the plane of paper. But they all have different linear(tangential) accelerations $a_{1}=r_{1} \alpha, a_{2}=r_{2} \alpha$ ...... , $a_{N}=r_{N} \alpha$. The force experienced by each particle is $f_{1}=m_{1} a_{1}=m_{1} r_{1} \alpha$, $f_{2}=m_{2} a_{2}=m_{2} r_{2} \alpha, \ldots . ., f_{N}=m_{N} a_{N}=m_{N} r_{N} \alpha$
These forces are tangential and their respective perpendicular distances are $r_{1}, r_{2}, \ldots, r_{N}$
Torque experienced is $\tau_{1}=f_{1} r_{1}=m_{1} r_{1}{ }^{2} \alpha, \tau_{2}=f_{2} r_{2}=m_{2} r_{2}{ }^{2} \alpha, \ldots \ldots . ., \tau_{N}=f_{N} r_{N}=m_{N} r_{N}{ }^{2} \alpha$ Magnitude of the resulting torque $\tau=\tau_{1}+\tau_{2}+\ldots \ldots+\tau_{N}$ $\tau=m_{1} r_{1}^{2} \alpha+m_{2} r_{2}^{2} \alpha+\ldots \ldots \ldots . .+m_{N} r_{N}^{2} \alpha=\left(m_{1} r_{1}^{2}+m_{2} r_{2}^{2}+\ldots \ldots \ldots . .+m_{N} r_{N}^{2}\right) \alpha$ $\tau=I \alpha$, where $I=m_{1} r_{1}{ }^{2}+m_{2} r_{2}{ }^{2}+\ldots .+m_{N} r_{N}{ }^{2}=\sum_{i=1}^{N} m_{i} r_{i}^{2}$ is the moment of Inertia (MI) of the object about the given axis of rotation.
NOTE: $\tau=I \alpha$ is analogous to $f=m a$ of linear motion, if moment of Inertia (I) replaces mass, which is its physical significance.

## Conservation of Angular Momentum:

Angular momentum $\bar{L}=\bar{r} X \bar{p}$ where $\bar{r}$ is the position vector from the axis of rotation and $\bar{p}$ is the linear momentum
Differentiating w.r.t. time, we get
$\frac{d \bar{L}}{d t}=\frac{d}{d t}(\bar{r} X \bar{p})=\bar{r} X \frac{d \bar{p}}{d t}+\frac{d \bar{r}}{d t} X \bar{p}$
$\frac{d \bar{L}}{d t}=\bar{r} X \bar{F}+m(\bar{v} X \bar{v})$, Since $\frac{d \bar{p}}{d t}=\bar{F} \quad$ and $\frac{d \bar{r}}{d t}=\bar{v}$
Since $\bar{v} X \bar{v}=\overline{0}$
Thus $\frac{d \bar{L}}{d t}=\bar{r} X \bar{F}$ which is moment of force of torque
$\tau=\frac{d \bar{L}}{d t} \ldots \ldots$ If $\tau=0, \frac{d \bar{L}}{d t}=0$, hence $\bar{L}=$ constant $($ conserved $)$
Hence in the absence of external unbalanced torque, $\bar{L}$ is conserved

## Examples

Ballet Dancers: During ice ballets the dancers come together (MI decreases, frequency increases) while taking rounds of small radius. While taking outer rounds the dancers outstretch their legs and arms. This will increase the MI and reduces the angular speed and linear speed. This is necessary to prevent slipping
Diving into a swimming pool during a competition: When standing on a diving board the diver stretches to have higher MI. Immediately on leaving the board the diver folds his body, to decrease MI, increase his frequency so he can complete more rounds in air. While just entering the pool diver stretches the body into a streamline shape for smooth entry into water.


C be the center of mass M Object mass
Axis of rotation MOP is parallel to axis ACB which is passing through center of mass and perpendicular to the plane.
$h$ is the distance between the parallel axes.
Consider a mass dm at D . The MI about the two axes is $I_{c}=\int(D C)^{2} d m$ and $I_{O}=\int(D O)^{2} d m$

$$
\begin{aligned}
& I_{O}=\int\left[(D N)^{2}+(O N)^{2}\right] d m \text { But ON }=O C+C N \\
& I_{O}=\int\left[(D N)^{2}+(O C+C N)^{2}\right] d m \\
&
\end{aligned}
$$

$$
I_{O}=\int\left[(D C)^{2}+h^{2}+2(h)(C N)\right] d m
$$

$$
=\int(D C)^{2} d m+h^{2} \int d m+2 h \int C N d m
$$

$$
\int(D C)^{2} d m=I_{c}, \quad \int d m=M, \quad \int C N d m=0
$$

since mass distribution is symmetric
Therfore, $I_{o}=I_{c}+M h^{2}$
The MI ( $I_{0}$ ) of an object about any axis is given by the sum of the moment of inertial about a parallel axis through the center of mass and product of the mass of the object and the square of the distance between the two axes ( $M h^{2}$ ).

## PERPENDICULAR AXIS Theorem:



Let us consider a rigid laminar object able to rotate about three mutually perpendicular axes $x, y$ and $z$. Let $x$ and $y$ axis be in the plane and $z$ axis perpendicular to the plane. All 3 axes meet at $O$. Let dm be the mass at point $P . x$ and $y$ are the perpendicular distance of $P$ from the $y$ and $x$ axis respectively. The distance from $Z$ axis will be $\sqrt{x^{2}+y^{2}}$. Let $I_{x}, I_{y}$ and $I_{z}$ be the MI of the body about the respective axes.

$$
\boldsymbol{I}_{\boldsymbol{Z}}=\int z^{2} d m=\int\left(x^{2}+y^{2}\right) d m=\int x^{2} d m+\int y^{2} d m=\boldsymbol{I}_{\boldsymbol{x}}+\boldsymbol{I}_{\boldsymbol{y}}
$$

The $\operatorname{MI}\left(I_{z}\right)$ of a laminar object about an axis (z) perpendicular to its plane is the sum of its MI about two mutually perpendicular axes ( $x$ and y) in its plane, all three axes being concurrent.

| Equation for translational motion | Analogous equation for rotational motion |
| :---: | :---: |
| $\mathrm{v}_{\text {ov }}=\frac{u+\mathrm{v}}{2}$ | $\omega_{a v}=\frac{\omega_{0}+\omega}{2}$ |
| $\begin{aligned} & a=\frac{d \mathrm{v}}{d t}=\frac{\mathrm{v}-u}{t} \\ & \therefore \mathrm{v}=u+a t \end{aligned}$ | $\begin{aligned} & \alpha=\frac{d \omega}{d t}=\frac{\omega-\omega_{0}}{t} \\ & \therefore \omega=\omega_{0}+\alpha t \end{aligned}$ |
| $\begin{aligned} & s=\mathrm{v}_{\mathrm{av}} \cdot t \\ & =\left(\frac{u+\mathrm{v}}{2}\right) t \\ & =u t+\frac{1}{2} a t^{2} \end{aligned}$ | $\begin{aligned} \theta & =\omega_{a v} \cdot t \\ & =\left(\frac{\omega_{0}+\omega}{2}\right) t \\ & =\omega_{0} t+\alpha t^{2} \end{aligned}$ |
| $\mathrm{v}^{2}=u^{2}+2 a s$ | $\omega^{2}=\omega_{0}^{2}+2 \alpha \theta$ |


| Translational motion |  | Rotational motion |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Quantity | Symbol/ <br> expression | Quantity | Symbol/ <br> expression | Inter-relation, <br> if possible |
| Linear <br> displacement | $\vec{s}$ | Angular <br> displacement | $\vec{\theta}$ | $\vec{s}=\vec{\theta} \times \vec{r}$ |
| Linear velocity | $\overrightarrow{\mathrm{v}}=\frac{d \vec{s}}{d t}$ | Angular velocity | $\vec{\omega}=\frac{d \vec{\theta}}{d t}$ | $\overrightarrow{\mathrm{v}}=\vec{\omega} \times \vec{r}$ |
| Linear <br> acceleration | $\vec{a}=\frac{d \overrightarrow{\mathrm{v}}}{d t}$ | Angular <br> acceleration | $\vec{\alpha}=\frac{d \vec{\omega}}{d t}$ | $\vec{\alpha}=\overrightarrow{\boldsymbol{a}} \times \vec{r}$ |
| Inertia or mass | $m$ | Rotational inertia <br> or moment of <br> inertia | $I$ | $I=\int r^{2} d m=\sum m_{i} r_{i}^{2}$ |
| Linear <br> momentum | $\vec{p}=m \overrightarrow{\mathrm{v}}$ | Angular <br> momentum | $\vec{L}=I \vec{\omega}$ | $\vec{L}=\vec{r} \times \vec{p}$ |
| Force | $\vec{f}=\frac{d \vec{p}}{d t}$ | Torque | $\vec{\tau}=\frac{d \vec{L}}{d t}$ | $\vec{\tau}=\vec{r} \times \vec{f}$ |
| Work | $W=\vec{f} \cdot \vec{s}$ | Work | $W=\vec{\tau} \cdot \vec{\theta}$ | ----- |
| Power | $P=\frac{d W}{d t}=\vec{f} \cdot \overrightarrow{\mathrm{v}}$ | Power | $P=\frac{d W}{d t}=\vec{\tau} \cdot \vec{\omega}$ | ------ |



Consider a disc of mass $M$ and negligible thickness and radius $R$. Let the axis of rotation pass through the center of the disc and perpendicular to it place. Let $\sigma=\frac{M}{A}=\frac{M}{\pi R^{2}}$ be uniform. A disc can be considered a sequence of rings. Consider one such ring at distance $r$ and mass dm and negligible thickness dr. Then $\sigma=\frac{d m}{2 \pi r d r}$. MI of this ring is $\mathrm{I}_{\mathrm{r}}=\mathrm{dm} . \mathrm{r}^{2}$
The MI of the disc will be got by integrating this from $r=0$ to $r=R$

$$
\begin{gathered}
I=\int_{o}^{R} d m r^{2}=\int_{o}^{R} \sigma 2 \pi r d r \cdot r^{2}=2 \pi \sigma \int_{o}^{R} r^{3} d r=2 \pi \sigma\left[\frac{r^{4}}{4}\right]_{0}^{R} \\
=2 \pi \frac{M}{\pi R^{2}} \frac{R^{4}}{4}=\frac{M R^{2}}{2}
\end{gathered}
$$

